量子开放系统理论与计算方法培训班 2017年5月22日~23日. 北京计算科学研究中心

开放系统量子力学 耗散子动力学(DEOM)理论

严以京

中国科学技术大学

数材: "Dissipation Equation of Motion Approach in Quantum Mechanics of Open Systems", Y. J. Yan, J. S. Jin, R. X. Xu, and X. Zheng, *Front. Phys.* **11**, 110306 (2016)

提要

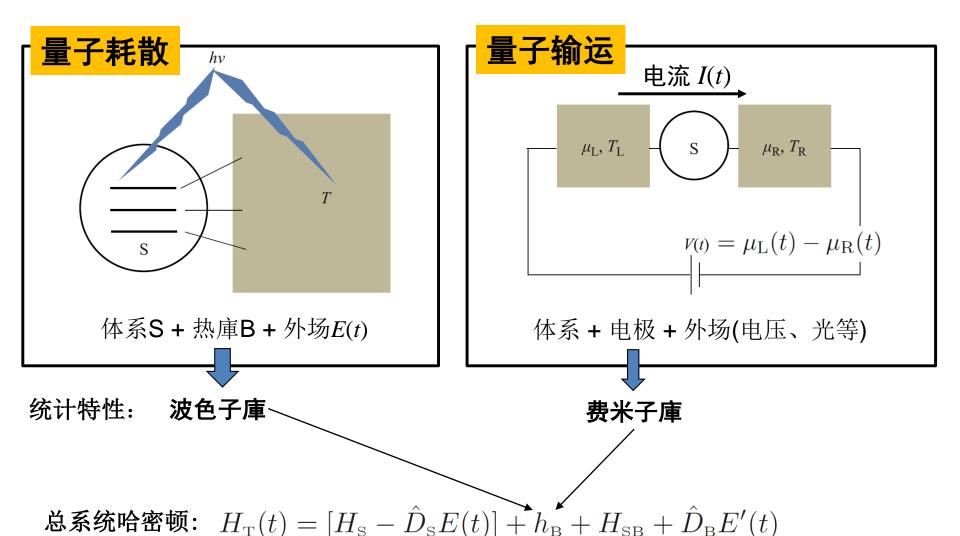
- **教材**(DEOM理论综述文章): "Dissipation Equation of Motion Approach in Quantum Mechanics of Open Systems", Y. J. Yan et al, *Front. Phys.* **11**, 110306 (2016)
- 授课ppt内容包含教材的§1-§3.41以及§4.1;课堂上会有些简单但重要的现场推导。 应用DEOM研究具体体系的信息将在ppt中给出,方便课后查阅和下载。
- DEOM理论是开放系统量子力学的一个最新发展,涵盖了多种已有的量子耗散动力学,如级联运动方程(HEOM)和主方程理论。费米型DEOM理论能精确、高效研究如量子输运、Kondo物理、Anderson杂质等问题。
- **DEOM**描述(**体系+杂化环境**)的**动力学、平衡态**,以及**非平衡态**的性质。传统量子力学的Hilbert空间方法,如薛定格、海森堡、相互作用图像等,都可以很方便的拓展到**DEOM**空间; **参见教材的§7**。
- 本节课程为(50+10)分钟,将重点讨论DEOM理论的物理思想、耗散子代数和定理、 以及耗散子级联运动方程的构造和方法。由于时间关系,本课程主要介绍波色型 DEOM理论,并适当兼顾讨论费米型理论的特点。

课前基础知识准备

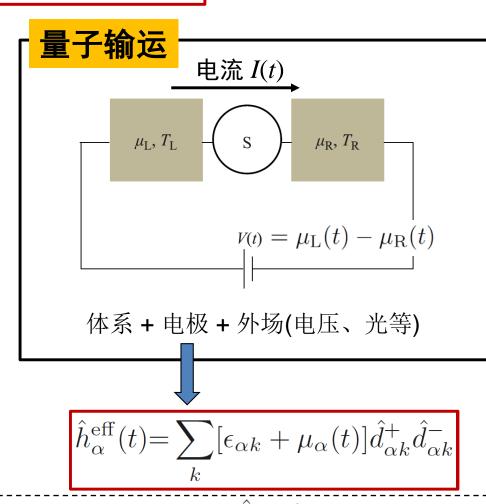
- 阅读教材,尤其是其中的§1-§2.3,§3.1,§4.1,§4.2
- 附上有关二级微扰主方程理论的一篇综述, Annu. Rev. Phys. Chem. **56**, 187-219 (2005), 自学其中的*§1-§3.1*, pp.187-194

内容

- 1 开放系统量子力学的研究对象和出发点
- 2 DEOM理论介绍和推导
- 3 总结以及相关方法发展简介



$$H_{\scriptscriptstyle \mathrm{T}}(t) = H(t) + \sum_{\alpha = \mathrm{L,R}} \hat{h}_{\alpha}^{\mathrm{eff}}(t) + H_{\scriptscriptstyle \mathrm{SB}}$$



总系统哈密顿: $H_{\text{T}}(t) = [H_{\text{B}} - \hat{D}_{\text{B}}E(t)] + h_{\text{B}} + H_{\text{BB}} + \hat{D}_{\text{B}}E'(t)$

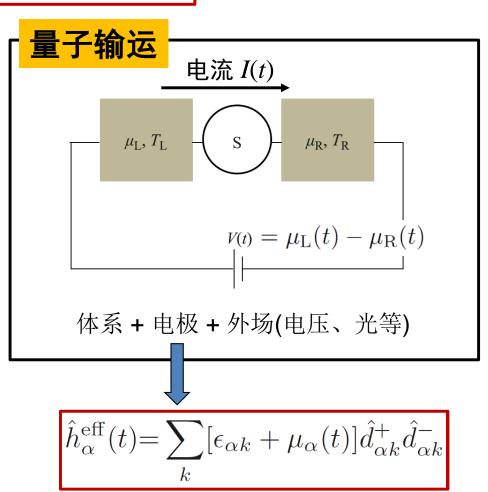
$$H_{\mathrm{T}}(t) = H(t) + \sum_{\alpha = \mathrm{L,R}} \hat{h}_{\alpha}^{\mathrm{eff}}(t) + H_{\mathrm{SB}}$$

$$H_{\rm SB} = \sum_{\alpha u} \left(\hat{a}_u^+ \hat{F}_{\alpha u}^- + \hat{F}_{\alpha u}^+ \hat{a}_u^- \right)$$

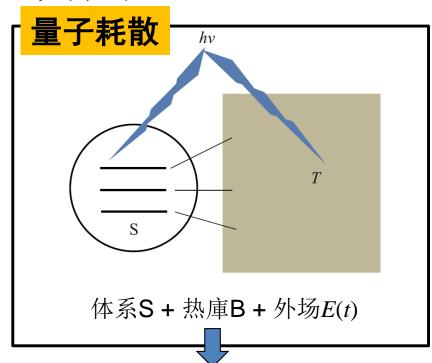
$$\hat{F}_{\alpha u}^{+} \equiv \sum_{k} t_{\alpha k u} \hat{d}_{\alpha k}^{+} = (\hat{F}_{\alpha u}^{-})^{\dagger}$$

电流算符

$$\hat{I}_{\alpha} \equiv -\frac{d}{dt}\hat{N}_{\alpha} = -i[H_{SB}, \hat{N}_{\alpha}]$$
$$= -i\sum_{u} \left(\hat{a}_{u}^{\dagger}\hat{F}_{\alpha u}^{-} - \hat{F}_{\alpha u}^{\dagger}\hat{a}_{u}^{-}\right)$$



本课程讨论

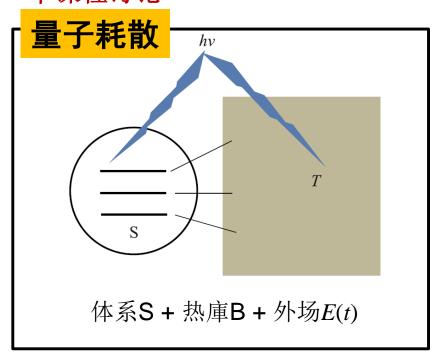


统计特性: 波色子庫

总系统哈密顿: $H_{\text{T}}(t) = [H_{\text{B}} - \hat{D}_{\text{B}}E(t)] + h_{\text{B}} + H_{\text{BB}} + \hat{D}_{\text{B}}E'(t)$

$$H_{\text{\tiny T}}(t) = H(t) + h_{\text{\tiny B}} + \sum_a [\hat{Q}_a^{\text{\tiny S}} - \zeta_a^{\text{\tiny B}} E(t)] \hat{F}_a^{\text{\tiny B}}$$

本课程讨论



波色子庫:
$$h_{\rm B} = \frac{1}{2} \sum_{j} \omega_{j} (p_{j}^{2} + x_{j}^{2})$$
 $\hat{F}_{a}^{\rm B} = \sum_{j} c_{aj} x_{j}$

总系统哈密顿: $H_{\text{T}}(t) = [H_{\text{S}} - \hat{D}_{\text{S}}E(t)] + h_{\text{B}} + H_{\text{SB}} + \hat{D}_{\text{B}}E'(t)$

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• 高斯环境影响的表征

杂化环境谱密度函数

$$h_{\rm B} = \frac{1}{2} \sum_{j} \omega_j (p_j^2 + x_j^2)$$

$$\hat{F}_a^{\rm B} = \sum_j c_{aj} x_j$$

$$H_{\text{\tiny T}}(t) = H(t) + h_{\text{\tiny B}} + \sum_a [\hat{Q}_a^{\text{\tiny S}} - \zeta_a^{\text{\tiny B}} E(t)] \hat{F}_a^{\text{\tiny B}}$$

• 高斯环境影响的表征

杂化环境谱密度函数

微观表达式 (for $\omega \geqslant 0$):

$$J_{ab}(\omega) = \frac{\pi}{2} \sum_{j} c_{aj} c_{bj} \delta(\omega - \omega_j)$$

热力学表达式 (波色子庫):

$$J_{ab}(\omega) = \frac{1}{2} \int_{-\infty}^{\infty} dt \, e^{i\omega t} \langle [\hat{F}_a^{\mathrm{B}}(t), \hat{F}_b^{\mathrm{B}}(0)] \rangle_{\mathrm{B}}$$

这里,
$$\hat{F}_a^{\mathrm{B}}(t) \equiv \mathrm{e}^{\mathrm{i}h_{\mathrm{B}}t/\hbar}\hat{F}_a^{\mathrm{B}}\mathrm{e}^{-\mathrm{i}h_{\mathrm{B}}t/\hbar}$$

$$\langle \hat{O} \rangle_{\rm B} \equiv \frac{\mathrm{tr_B}(\hat{O}e^{-\beta h_{\rm B}})}{\mathrm{tr_B}e^{-\beta h_{\rm B}}}$$

$$h_{\rm B} = \frac{1}{2} \sum_{j} \omega_{j} (p_{j}^{2} + x_{j}^{2})$$

$$\hat{F}_a^{\mathrm{B}} = \sum_j c_{aj} x_j$$

$$H_{\text{\tiny T}}(t) = H(t) + h_{\text{\tiny B}} + \sum_{a} [\hat{Q}_{a}^{\text{\tiny S}} - \zeta_{a}^{\text{\tiny B}} E(t)] \hat{F}_{a}^{\text{\tiny B}}$$

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$$J_{ab}^*(\omega) = -J_{ab}(-\omega) = J_{ba}(\omega)$$

$$h_{\rm B} = \frac{1}{2} \sum_{j} \omega_j (p_j^2 + x_j^2)$$

$$\hat{F}_a^{\mathrm{B}} = \sum_j c_{aj} x_j$$

$$H_{\text{\tiny T}}(t) = H(t) + h_{\text{\tiny B}} + \sum_a [\hat{Q}_a^{\text{\tiny S}} - \zeta_a^{\text{\tiny B}} E(t)] \hat{F}_a^{\text{\tiny B}}$$

• 高斯环境影响的表征

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热力学表达式 (波色子庫):

$$J_{ab}(\omega) = \frac{1}{2} \int_{-\infty}^{\infty} dt \, e^{i\omega t} \langle [\hat{F}_a^{\mathrm{B}}(t), \hat{F}_b^{\mathrm{B}}(0)] \rangle_{\mathrm{B}}$$

杂化环境关联函数,以及涨落耗散定理

$$\langle \hat{F}_a^{\mathrm{B}}(t)\hat{F}_b^{\mathrm{B}}(0)\rangle_{\mathrm{B}} = \frac{1}{\pi} \int_{-\infty}^{\infty} \mathrm{d}\omega \,\mathrm{e}^{-\mathrm{i}\omega t} \frac{J_{ab}(\omega)}{1 - \mathrm{e}^{-\beta\omega}}$$

$$h_{\rm B} = \frac{1}{2} \sum_{j} \omega_j (p_j^2 + x_j^2)$$

$$\hat{F}_a^{\mathrm{B}} = \sum_j c_{aj} x_j$$

$$H_{\rm T}(t) = H(t) + h_{\rm B} + \sum_a [\hat{Q}_a^{\rm S} - \zeta_a^{\rm B} E(t)] \hat{F}_a^{\rm B}$$

任意

定理:

杂化环境关联函数完全确定了高斯 环境对开放系统动力学的影响

$$h_{\rm B} = \frac{1}{2} \sum_{j} \omega_j (p_j^2 + x_j^2)$$

$$\hat{F}_a^{\mathrm{B}} = \sum_j c_{aj} x_j$$

$$\langle \hat{F}_a^{\mathrm{B}}(t)\hat{F}_b^{\mathrm{B}}(0)\rangle_{\mathrm{B}} = \frac{1}{\pi} \int_{-\infty}^{\infty} \mathrm{d}\omega \,\mathrm{e}^{-\mathrm{i}\omega t} \frac{J_{ab}(\omega)}{1 - \mathrm{e}^{-\beta\omega}}$$

$$H_{\mathrm{T}}(t) = H(t) + h_{\mathrm{B}} + \sum_{a} [\hat{Q}_{a}^{\mathrm{S}} - \zeta_{a}^{\mathrm{B}} E(t)] \hat{F}_{a}^{\mathrm{B}}$$

波色谱密度对称性: $J_{ab}^*(\omega) = -J_{ab}(-\omega) = J_{ba}(\omega)$

$$\langle \hat{F}_a^{\mathrm{B}}(t)\hat{F}_b^{\mathrm{B}}(0)\rangle_{\mathrm{B}} = \frac{1}{\pi} \int_{-\infty}^{\infty} \mathrm{d}\omega \,\mathrm{e}^{-\mathrm{i}\omega t} \frac{J_{ab}(\omega)}{1 - \mathrm{e}^{-\beta\omega}}$$

逆时序对称性: $\langle \hat{F}_b^{\mathrm{B}}(0)\hat{F}_a^{\mathrm{B}}(t)\rangle_{\mathrm{B}}=\langle \hat{F}_a^{\mathrm{B}}(t)\hat{F}_b^{\mathrm{B}}(0)\rangle_{\mathrm{B}}^*$

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- 1 开放系统量子力学的研究对象和出发点
- 2 DEOM理论介绍和推导
- 3 总结以及相关方法发展简介

2.1 DEOM理论构建的出发点: 耗散子的引入

$$H_{\mathrm{T}}(t) = H(t) + h_{\mathrm{B}} + \sum_{a} [\hat{Q}_{a}^{\mathrm{S}} - \zeta_{a}^{\mathrm{B}} E(t)] \hat{F}_{a}^{\mathrm{B}}$$

波色谱密度对称性: $J_{ab}^*(\omega)=-J_{ab}(-\omega)=J_{ba}(\omega)$

$$\langle \hat{F}_a^{\mathrm{B}}(t)\hat{F}_b^{\mathrm{B}}(0)\rangle_{\mathrm{B}} = \frac{1}{\pi} \int_{-\infty}^{\infty} \mathrm{d}\omega \,\mathrm{e}^{-\mathrm{i}\omega t} \frac{J_{ab}(\omega)}{1 - \mathrm{e}^{-\beta\omega}} = \sum_{k=1}^{K} \eta_{abk} \mathrm{e}^{-\gamma_k t}$$

逆时序对称性:
$$\langle \hat{F}_b^{\mathrm{B}}(0)\hat{F}_a^{\mathrm{B}}(t)\rangle_{\mathrm{B}} = \langle \hat{F}_a^{\mathrm{B}}(t)\hat{F}_b^{\mathrm{B}}(0)\rangle_{\mathrm{B}}^* = \sum_{k=1}^{N} \eta_{ab\bar{k}}^* \mathrm{e}^{-\gamma_k t}$$

下标
$$\bar{k} \in \{k=1,\cdots,K\}$$
 由 $\gamma_{\bar{k}} \equiv \gamma_k^*$ 而定义

2.2 DEOM理论构建的出发点: 耗散子的引入

$$H_{\text{\tiny T}}(t) = H(t) + h_{\text{\tiny B}} + \sum_{a} [\hat{Q}_{a}^{\text{\tiny S}} - \zeta_{a}^{\text{\tiny B}} E(t)] \hat{F}_{a}^{\text{\tiny B}}$$

$$\left\langle \hat{F}_{a}^{\mathrm{B}}(t)\hat{F}_{b}^{\mathrm{B}}(0)\right\rangle_{\mathrm{B}} = \sum_{k=1}^{K} \eta_{abk} \mathrm{e}^{-\gamma_{k}t}$$
 和 $\left\langle \hat{F}_{b}^{\mathrm{B}}(0)\hat{F}_{a}^{\mathrm{B}}(t)\right\rangle_{\mathrm{B}} = \sum_{k=1}^{K} \eta_{ab\bar{k}}^{*} \mathrm{e}^{-\gamma_{k}t}$

2.2 DEOM理论构建的出发点: 耗散子的引入

$$H_{\text{\tiny T}}(t) = H(t) + h_{\text{\tiny B}} + \sum_a [\hat{Q}_a^{\text{\tiny S}} - \zeta_a^{\text{\tiny B}} E(t)] \hat{F}_a^{\text{\tiny B}}$$

$$\left\langle \left\langle \hat{F}_{a}^{\mathrm{B}}(t)\hat{F}_{b}^{\mathrm{B}}(0)\right\rangle _{\mathrm{B}}=\sum_{k=1}^{K}\eta _{abk}\mathrm{e}^{-\gamma _{k}t}\right\rangle$$

$$\left\langle \left\langle \hat{F}_{a}^{\mathrm{B}}(t)\hat{F}_{b}^{\mathrm{B}}(0)\right\rangle _{\mathrm{B}}=\sum_{k=1}^{K}\eta_{abk}\mathrm{e}^{-\gamma_{k}t}\quad \mathbf{A}\quad \left\langle \hat{F}_{b}^{\mathrm{B}}(0)\hat{F}_{a}^{\mathrm{B}}(t)\right\rangle _{\mathrm{B}}=\sum_{k=1}^{K}\eta_{ab\bar{k}}^{*}\mathrm{e}^{-\gamma_{k}t}$$



耗散子(杂化环境的统计等效准粒子),表示为 $\left|\hat{F}_{a}^{\mathrm{B}} ight| \sum \hat{f}_{ak}$,满足

$$\hat{F}_a^{ ext{ iny B}} \equiv \sum_{k=1}^K \hat{f}_{ak}$$
,满足

$$\langle \hat{f}_{ak}(t)\hat{f}_{bj}(0)\rangle_{\rm B} = \delta_{kj}\eta_{abk}e^{-\gamma_k t}$$

和
$$\langle \hat{f}_{bj}(0)\hat{f}_{ak}(t)\rangle_{_{\mathrm{B}}} = \delta_{kj}\eta_{ab\bar{k}}^*\mathrm{e}^{-\gamma_k t}$$

2.2 DEOM理论构建的出发点: 耗散子的引入

$$H_{\rm T}(t) = H(t) + h_{\rm B} + \sum_{ak} [\hat{Q}_a^{\rm S} - \zeta_a^{\rm B} E(t)] \hat{f}_{ak}$$

$$\hat{F}_a^{\mathrm{B}} \equiv \sum_{k=1}^K \hat{f}_{ak}$$

$$\left\langle \hat{f}_{ak}(t)\hat{f}_{bj}(0)\right\rangle_{\mathrm{B}} = \delta_{kj}\eta_{abk}\mathrm{e}^{-\gamma_k t}$$
 \mathbf{a} $\left\langle \hat{f}_{bj}(0)\hat{f}_{ak}(t)\right\rangle_{\mathrm{B}} = \delta_{kj}\eta_{ab\bar{k}}^*\mathrm{e}^{-\gamma_k t}$

$$\langle \hat{f}_{bj}(0)\hat{f}_{ak}(t)\rangle_{\rm B} = \delta_{kj}\eta_{ab\bar{k}}^* e^{-\gamma_k t}$$

耗散子代数:
$$\operatorname{tr}_{\mathrm{B}}\left[\left(\frac{\partial}{\partial t}\hat{f}_{ak}\right)_{\mathrm{B}}\rho_{\mathrm{T}}(t)\right] = -\gamma_{k}\operatorname{tr}_{\mathrm{B}}\left[\hat{f}_{ak}\rho_{\mathrm{T}}(t)\right]$$

2.3 多体耗散子密度算符(DDO)和代数

$$H_{\mathrm{T}}(t) = H(t) + h_{\mathrm{B}} + \sum_{ak} [\hat{Q}_{a}^{\mathrm{S}} - \zeta_{a}^{\mathrm{B}} E(t)] \hat{f}_{ak}$$

$$\rho_{\mathbf{n}}^{(n)}(t) \equiv \operatorname{tr}_{\mathrm{B}} \left[\left(\prod_{ak} \hat{f}_{ak}^{n_{ak}} \right)^{\circ} \rho_{\mathrm{T}}(t) \right]$$

(n体DDO)

- ho $n \equiv \{n_{ak}\}$; with $n_{ak} = 0, 1, 2, \cdots$ (构型)
- $\rightarrow n = \sum n_{ak}$
- ightharpoonup $(\cdots)^{\circ}$ 不可约符号; 波色型耗散子满足 $(\hat{f}_k\hat{f}_j)^{\circ}=(\hat{f}_j\hat{f}_k)^{\circ}$

$$\langle \hat{f}_{ak}(t)\hat{f}_{bj}(0)\rangle_{\rm B} = \delta_{kj}\eta_{abk}e^{-\gamma_k t}$$

$$\left\langle \hat{f}_{ak}(t)\hat{f}_{bj}(0)\right\rangle_{\mathrm{B}} = \delta_{kj}\eta_{abk}\mathrm{e}^{-\gamma_k t} \qquad \mathbf{\pi} \qquad \left\langle \hat{f}_{bj}(0)\hat{f}_{ak}(t)\right\rangle_{\mathrm{B}} = \delta_{kj}\eta_{ab\bar{k}}^*\mathrm{e}^{-\gamma_k t}$$

耗散子代数:

$$\operatorname{tr}_{\mathrm{B}}\left[\left(\frac{\partial}{\partial t}\hat{f}_{ak}\right)_{\mathrm{B}}\rho_{\mathrm{T}}(t)\right] = -\gamma_{k}\operatorname{tr}_{\mathrm{B}}\left[\hat{f}_{ak}\rho_{\mathrm{T}}(t)\right]$$

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$$H_{\mathrm{T}}(t) = H(t) + h_{\mathrm{B}} + \sum_{ak} [\hat{Q}_{a}^{\mathrm{S}} - \zeta_{a}^{\mathrm{B}} E(t)] \hat{f}_{ak}$$

$$\rho_{\boldsymbol{n}}^{(n)}(t) \equiv \operatorname{tr}_{\mathrm{B}} \left[\left(\prod_{ak} \hat{f}_{ak}^{n_{ak}} \right)^{\circ} \rho_{\mathrm{T}}(t) \right]$$

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- ho $n \equiv \{n_{ak}\}$; with $n_{ak} = 0, 1, 2, \cdots$ (构型)
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$$\langle \hat{f}_{ak}(t)\hat{f}_{bj}(0)\rangle_{\mathbf{p}} = \delta_{kj}\eta_{abk}e^{-\gamma_k t}$$

$$\left\langle \hat{f}_{ak}(t)\hat{f}_{bj}(0)\right\rangle_{\mathrm{B}} = \delta_{kj}\eta_{abk}\mathrm{e}^{-\gamma_k t} \qquad \mathbf{\pi} \qquad \left\langle \hat{f}_{bj}(0)\hat{f}_{ak}(t)\right\rangle_{\mathrm{B}} = \delta_{kj}\eta_{ab\bar{k}}^*\mathrm{e}^{-\gamma_k t}$$

耗散子代数:

$$\operatorname{tr}_{\mathrm{B}}\left[\left(\frac{\partial}{\partial t}\hat{f}_{ak}\right)_{\mathrm{B}}\rho_{\mathrm{T}}(t)\right] = -\gamma_{k}\operatorname{tr}_{\mathrm{B}}\left[\hat{f}_{ak}\rho_{\mathrm{T}}(t)\right]$$

$$\left[i \operatorname{tr}_{\mathbf{B}} \left\{ \left(\prod_{ak} \hat{f}_{ak}^{n_{ak}} \right)^{\circ} \left[h_{\mathbf{B}}, \rho_{\mathbf{T}} \right] \right\} = \left(\sum_{ak} n_{ak} \gamma_k \right) \rho_{\mathbf{n}}^{(n)} \right]$$

2.3 多体耗散子密度算符(DDO)和代数

$$H_{\mathrm{T}}(t) = H(t) + h_{\mathrm{B}} + \sum_{ak} [\hat{Q}_{a}^{\mathrm{S}} - \zeta_{a}^{\mathrm{B}} E(t)] \hat{f}_{ak}$$

$$\rho_{\mathbf{n}}^{(n)}(t) \equiv \operatorname{tr}_{\mathrm{B}} \left[\left(\prod_{ak} \hat{f}_{ak}^{n_{ak}} \right)^{\circ} \rho_{\mathrm{T}}(t) \right]$$

(n体DDO)

- $\boldsymbol{n} \equiv \{n_{ak}\}$; with $n_{ak} = 0, 1, 2, \cdots$ (构型)
- $n = \sum n_{ak}$
- \blacktriangleright $(\cdots)^{\circ}$ 不可约符号, 波色型耗散子满足 $(\hat{f}_{ak}\hat{f}_{bj})^{\circ} = (\hat{f}_{bj}\hat{f}_{ak})^{\circ}$

$$\langle \hat{f}_{ak}(t)\hat{f}_{bj}(0)\rangle_{_{\mathbf{R}}} = \delta_{kj}\eta_{abk}e^{-\gamma_k t}$$

$$\left\langle \hat{f}_{ak}(t)\hat{f}_{bj}(0)\right\rangle_{\mathrm{B}} = \delta_{kj}\eta_{abk}\mathrm{e}^{-\gamma_k t} \quad \mathbf{m} \quad \left\langle \hat{f}_{bj}(0)\hat{f}_{ak}(t)\right\rangle_{\mathrm{B}} = \delta_{kj}\eta_{ab\bar{k}}^*\mathrm{e}^{-\gamma_k t}$$

$$\operatorname{tr}_{\mathrm{B}} \left[\left(\prod_{ak} \hat{f}_{ak}^{n_{ak}} \right)^{\circ} \hat{f}_{bj} \rho_{\mathrm{T}}(t) \right] = \sum_{ak} n_{ak} \left\langle \hat{f}_{ak} \hat{f}_{bj} \right\rangle_{\mathrm{B}}^{>} \rho_{\boldsymbol{n}_{ak}}^{(n-1)}(t) + \rho_{\boldsymbol{n}_{bj}}^{(n+1)}(t)$$

定理
$$\operatorname{tr}_{\mathrm{B}}\left[\left(\prod_{ak} \hat{f}_{ak}^{n_{ak}}\right)^{\circ} \rho_{\mathrm{T}}(t) \hat{f}_{bj}\right] = \sum_{ak} n_{ak} \langle \hat{f}_{bj} \hat{f}_{ak} \rangle_{\mathrm{B}}^{<} \rho_{\boldsymbol{n}_{ak}}^{(n-1)}(t) + \rho_{\boldsymbol{n}_{bj}}^{(n+1)}(t)$$

2.4 DEOM/HEOM的代数推导(课堂)

$$H_{\mathrm{T}}(t) = H(t) + h_{\mathrm{B}} + \sum_{ak} [\hat{Q}_{a}^{\mathrm{S}} - \zeta_{a}^{\mathrm{B}} E(t)] \hat{f}_{ak}$$

$$\rho_{\mathbf{n}}^{(n)}(t) \equiv \operatorname{tr}_{\mathrm{B}} \left[\left(\prod_{ak} \hat{f}_{ak}^{n_{ak}} \right)^{\circ} \rho_{\mathrm{T}}(t) \right]$$

级联方程(HEOM)

(推导: 从薛定格-刘维尔方程出发,利用耗散子代数完成)

$$\dot{\rho}_{\mathbf{n}}^{(n)} = -\left(i\mathcal{L}(t) + \sum_{ak} n_{ak}\gamma_k\right)\rho_{\mathbf{n}}^{(n)} - i\sum_{ak}\left(n_{ak}\mathcal{C}_{ak}^{\text{eff}}(t)\rho_{\mathbf{n}_{ak}^-}^{(n-1)} + \mathcal{A}_a\rho_{\mathbf{n}_{ak}^+}^{(n+1)}\right)$$

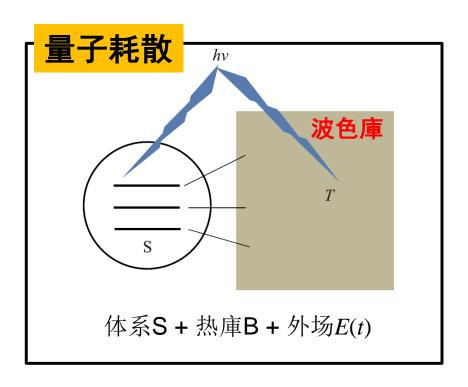
$$C_{ak}^{\text{eff}}(t) \equiv C_{ak} - E(t) \sum_{b} (\eta_{abk} - \eta_{ab\bar{k}}^*) \zeta_b^{\text{B}}, \qquad \mathcal{A}_a \hat{O} \equiv \left[\hat{Q}_a^{\text{S}}, \hat{O} \right]$$
$$C_{ak} \hat{O} \equiv \sum_{b} \left(\eta_{abk} \hat{Q}_b^{\text{S}} \hat{O} - \eta_{ab\bar{k}}^* \hat{O} \hat{Q}_b^{\text{S}} \right)$$

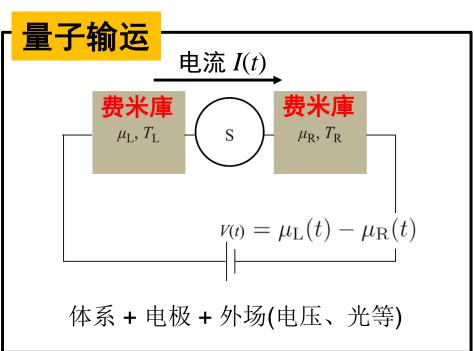
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3.1 DEOM理论的普适性

无穷自由度环境的效应 ⇒ 若干个耗散子的作用





耗散子: 统计独立准粒子

$$(\hat{f}_k \hat{f}_j)^{\circ} = (\hat{f}_j \hat{f}_k)^{\circ}$$
 (bosonic); $(\hat{f}_k \hat{f}_j)^{\circ} = -(\hat{f}_j \hat{f}_k)^{\circ}$ (fermionic).

3.2 DEOM理论 = HEOM + 耗散子代数

HEOM

$$\dot{\rho}_{\bm{n}}^{(n)} = -\left[\mathrm{i}\mathcal{L}(t) + \gamma_{\bm{n}}^{(n)}\right]\rho_{\bm{n}}^{(n)} + \left\{\rho_{\bm{n}^{-}}^{(n-1)}\right\} + \left\{\rho_{\bm{n}^{+}}^{(n+1)}\right\}$$

多体DDO

$$\rho_{\mathbf{n}}^{(n)}(t) \equiv \operatorname{tr}_{\mathrm{B}} \left[\left(\prod_{ak} \hat{f}_{ak}^{n_{ak}} \right)^{\circ} \rho_{\mathrm{T}}(t) \right]$$

波色型
$$(\hat{f}_{ak}\hat{f}_{bj})^{\circ}=(\hat{f}_{bj}\hat{f}_{ak})^{\circ}$$

耗散子: 统计独立准粒子

$$\begin{pmatrix} \hat{f}_{ak}(t)\hat{f}_{bj}(0) \rangle_{\mathrm{B}} = \delta_{kj}\eta_{abk}\mathrm{e}^{-\gamma_k t} \\ \langle \hat{f}_{bj}(0)\hat{f}_{ak}(t) \rangle_{\mathrm{B}} = \delta_{kj}\eta_{ab\bar{k}}^*\mathrm{e}^{-\gamma_k t} \end{pmatrix}$$

广义Wick's定理

$$\operatorname{tr}_{\mathrm{B}}\left[\left(\prod_{ak} \hat{f}_{ak}^{n_{ak}}\right)^{\circ} \hat{f}_{bj} \rho_{\mathrm{T}}(t)\right] = \sum_{ak} n_{ak} \langle \hat{f}_{ak} \hat{f}_{bj} \rangle_{\mathrm{B}}^{>} \rho_{\boldsymbol{n}_{ak}}^{(n-1)}(t) + \rho_{\boldsymbol{n}_{bj}}^{(n+1)}(t)$$

$$\langle \hat{f}_{ak} \hat{f}_{bj} \rangle_{\rm B}^{>} \equiv \langle \hat{f}_{ak} (0+) \hat{f}_{bj} \rangle = \delta_{kj} \eta_{abk}$$

3.3 DEOM / HEOM 方法的发展

- ▶ 波色型HEOM最早文章为 教材中的参考文献【12】; 费米型的为【18】。 (HEOM等价于Feynman-Venon影响泛函的路径积分公式,其动力学变量,除了体系约化密度矩阵以外,仅为数学上无物理意义的辅助量)。基于HEOM、严格的化学动力学速率公式可参见 J. Phys. Chem. A 120, 3241 (2016)
- ➤ **波色型HEOM**的程序包可参见日本京都大学化学系**Tanimura**教授的网页。 国内有许多研究小组也都有自己的**HEOM**程序
- ▶ **费米型HEOM**的程序包可参见中科大郑晓教授的网页。结合第一性原理 (DFT+U)计算,HEOM-QUICK程序包己广泛用于强关联电子体系及其 输运特性的研究 [参见综述 "HEOM-QUICK: A program for accurate, efficient and universal characterization of strongly correlated quantum impurity systems", L. Z Ye, X. L Wang, D. Hou, R. X Xu, X. Zheng, and Y. J. Yan, *WIREs Comp. Mol. Sci.* 6, 608-638 (2016)]
- ▶ DEOM理论拓展了HEOM方法,使得人们可以研究体系和环境的相干/纠缠动力学,以及平衡或非平衡稳态特性;参加教材中§7。体系和环境的相干动力学可以在实验上探测 [相关研究参见教材中的【22,23】,以及Science China Chem. 58 (12), 1816 (2015); Chem. Phys. 481, 237 (2016); J. Chem. Phys. 145, 204109 (2016)]

Thanks